

Upper limits for a narrow resonance in the reaction $p + p \rightarrow K^+ + (\Lambda p)$

A. Budzanowski*,⁹ A. Chatterjee,¹¹ H. Clement,¹³ P. Hawranek,¹⁰ F. Hinterberger,^{2,†} R. Jahn,² R. Joosten,² K. Kilian,^{5,6} Da. Kirillov,^{4,5} Di. Kirillov,³ S. Kliczewski,⁹ D. Kolev,¹² M. Kravcikova,⁷ M. Lesiak,¹⁰ H. Machner,⁴ A. Magiera,¹⁰ G. Martinska,⁸ N. Piskunov,³ D. Protić,⁵ J. Ritman,^{5,6} P. von Rossen,^{5,6} J. Roy,¹¹ A. Sibirtsev,¹ I. Sitnik,³ R. Siudak,⁹ R. Tsenov,¹² K. Ulbrich,² J. Urban,⁸ and G. J. Wagner¹³

(The HIRES Collaboration)

¹*CSSM, School of Chemistry and Physics, University of Adelaide, Adelaide SA 5005, Australia*

²*Helmholtz-Institut für Strahlen- und Kernphysik der Universität Bonn, 53115 Bonn, Germany*

³*Laboratory for High Energies, JINR Dubna, Russia*

⁴*Fachbereich Physik, Universität Duisburg-Essen, Duisburg, Germany*

⁵*Institut für Kernphysik, Forschungszentrum Jülich, Jülich, Germany*

⁶*Jülich Center for Hadron Physics, Forschungszentrum Jülich, Jülich, Germany*

⁷*Technical University, Kosice, Kosice, Slovakia*

⁸*P. J. Safarik University, Kosice, Slovakia*

⁹*Institute of Nuclear Physics, PAN, Krakow, Poland*

¹⁰*Institute of Physics, Jagellonian University, Krakow, Poland*

¹¹*Nuclear Physics Division, BARC, Mumbai, India*

¹²*Physics Faculty, University of Sofia, Sofia, Bulgaria*

¹³*Physikalisches Institut, Universität Tübingen, Tübingen, Germany*

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The reaction $pp \rightarrow K^+ + (\Lambda p)$ has been measured at $T_p = 1.953$ GeV and $\Theta = 0^\circ$ with a high missing-mass resolution in order to study the Λp final state interaction. Narrow $S = -1$ resonances predicted by bag model calculations are not visible in the missing-mass spectrum. Small structures observed in a previous experiment are not confirmed. Upper limits for the production cross section of a narrow resonance are deduced for missing-masses between 2058 and 2105 MeV.

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I. INTRODUCTION

A high-resolution study of the reaction $p + p \rightarrow K^+ + (\Lambda p)$ has been performed by the HIRES Collaboration [1, 2] using the proton beam of the Cooler Synchrotron COSY [3] and the magnetic spectrograph BIG KARL [4, 5] at the Research Center Jülich. The aim of the experiment was to study the Λp final state interaction (FSI) and to search for narrow strangeness $S = -1$ resonances. Concerning the FSI, first results have been published [1]. The present paper deals with the search for a narrow strangeness $S = -1$ resonance. We use only the missing-mass data measured at 1.953 GeV beam energy [1] and ignore a small portion of data below the ΣN threshold taken at some higher beam energy [2].

Predictions of strange dibaryons are summarized in a recent review by Gal [6]. A confined six-quark state $(Q^6)_1$ with $S = -2$, the so called H-dibaryon (from Hexaquark), has been predicted to be the lowest-lying dibaryon state [7]. It has been of prime interest both theoretically and experimentally. In the strangeness -1 sector narrow dibaryon resonances D_s and D_t are pre-

dicted to be located at about 55 MeV and 95 MeV above the Λp threshold [8–11]. The width of the lowest resonance D_s is estimated to be less than about 1 MeV.

Experimentally, the inclusive reaction $pp \rightarrow K^+ + (\Lambda p)$ was first studied with a rather low missing-mass resolution [12–14]. There were also exclusive measurements of the reaction $pp \rightarrow K^+ \Lambda p$ by the COSY-TOF Collaboration [15–17] and total cross section measurements by the COSY-11 [18] and COSY-ANKE Collaboration [19, 20]. The first high-resolution measurement of $pp \rightarrow K^+ + (\Lambda p)$ has been performed at SATURNE II with proton beam energies of 2.3 and 2.7 GeV using the spectrometer SPES4 [21]. The outgoing kaons were detected at forward angles with a high momentum resolution in the focal plane of the SPES4 spectrometer. The missing-mass spectra show characteristic enhancements near the Λp and ΣN thresholds. At 2.3 GeV and $\Theta = 10^\circ$, a sharp peak has been observed in the missing-mass spectrum at 2096.5 ± 1.5 MeV above the background which is due to the Λp continuum. The peak amounts to about 3.5 standard deviations. A similar peak corresponding to a missing-mass of 2098.0 ± 1.5 MeV has been observed at 2.7 GeV beam energy and 12.6° . The statistical accuracies were not high enough to exclude an accidental statistical fluctuation, however, the peak has been observed twice under different experimental and kinematical conditions and the peak energies coincide within the

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†corresponding author Frank Hinterberger: fh@hiskp.uni-bonn.de

experimental error. It is also interesting to note that the observed peaks were located near the predicted mass of the lowest $S = -1$ dibaryon D_s [8, 9]. Therefore, the reaction $pp \rightarrow K^+ + (\Lambda p)$ has been studied with an especially high accuracy in the missing-mass range between 2090 and 2110 MeV.

In Sec. II we give a short description of the experiment. In Sec. III we sketch the predictions of narrow dibaryon resonances. The effect of the lowest dibaryon D_s in the total cross section of the free $\Lambda p \rightarrow \Lambda p$ scattering is discussed in Sec. IV. The formalism to describe a narrow resonance embedded in the continuum of the reaction $pp \rightarrow K^+ + (\Lambda p)$ is presented in Sec. V. In Sec. VI we deduce upper limits for the production cross section of the predicted resonance D_s . Conclusion and discussion follow in Sec. VII.

II. EXPERIMENT

Here, we give a short description of the experiment which was already reported in [1]. The reaction $p + p \rightarrow K^+ + (\Lambda p)$ was measured at 0° using the proton beam from the cooler synchrotron COSY, the magnetic spectrograph BIG KARL [5] and a 1.0 cm thick liquid hydrogen target (see Fig. 1). The momentum of the incoming proton beam was 2.735 GeV/c corresponding to a kinetic energy of 1.953 GeV. The scattered particles in the momentum range 930 - 1110 MeV/c were detected in the focal plane using two stacks of multiwire drift chambers, two threshold Cherenkov detectors [22], and two scintillator hodoscopes. The absolute beam momentum was found by measuring simultaneously the tracks of K^+ particles and deuterons from the reactions $p + p \rightarrow K^+ + (\Lambda p)$ and $p + p \rightarrow d + \pi^+$ at fixed BIG KARL momentum 1070 MeV/c and fitting the kinematic parameters. The absolute precision of the beam momentum was 0.15 MeV/c. The ratio of beam momentum to scattered particle momentum allowed for a measurement at 0° . In the first dipole magnet of BIG KARL the beam was magnetically separated from the scattered particles and guided through the side exit of the outer yoke into the beam dump. Thus, the huge background from dumping the beam within the spectrometer was avoided. Particle identification was performed using the energy loss (ΔE) and time of flight (TOF) information from the scintillator hodoscopes. In addition two threshold Cherenkov detectors [22] were used in order to achieve a very high pion suppression factor of 10^5 . The momentum of the kaon was measured and the missing-mass of the Λp system was deduced. In order to cover the missing-mass range 2050 - 2110 MeV the data were taken using three overlapping settings of the spectrograph (mean momenta: 1070, 1010 and 960 MeV/c). The incoming beam was not changed during those measurements. The relative precision of the momenta of the three settings was 0.1 MeV/c.

Acceptance corrections with respect to solid angle and momentum were taken from Monte Carlo calcula-

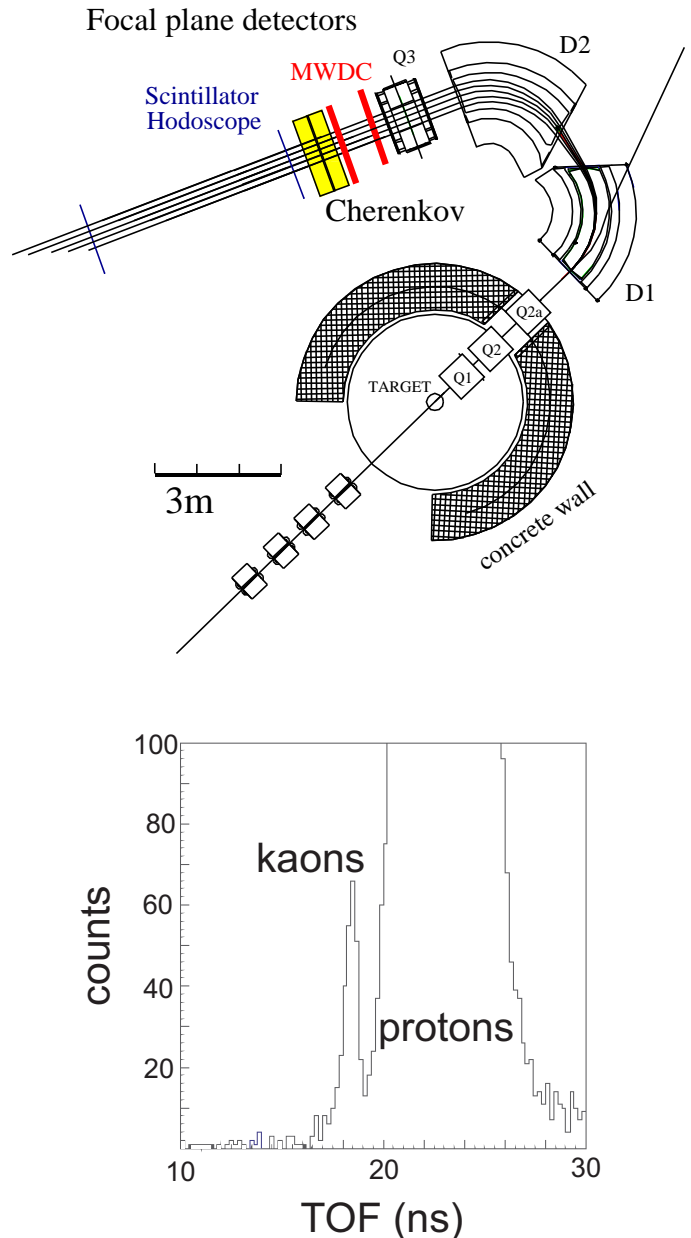


FIG. 1: Top: Layout of the magnetic spectrograph BIG KARL. The charged particle tracks are detected in the focal plane using two stacks of multiwire drift chambers, two threshold Cherenkov detectors, and two scintillator hodoscopes. Bottom: TOF spectrum with pion suppression by two Cherenkov detectors.

tions. The acceptance correction functions contained the magnetic spectrograph momentum acceptance around 0° with cuts on the horizontal emission angle $\Theta_x = \arctan(p_x/p_z)$ and the vertical emission angle $\Theta_y = \arctan(p_y/p_z)$ with p_x and p_y transversal momentum components and p_z the longitudinal component. The cuts for a given missing-mass bin corresponded to the measured solid angle $d\Omega$. They contained also the de-

detector efficiency corrections. The detector efficiency included efficiency of scintillator detectors which determined trigger and particle identification and magnetic spectrograph efficiency with field inhomogeneity at the edges of the acceptance. Acceptance corrections as determined by Monte Carlo simulations were checked by use of experimental distributions of simultaneously measured pions. They are shown in Fig. 2. The acceptance correction function of the 960 MeV/ c setting looks different due to a slightly different Θ_x and Θ_y cut.

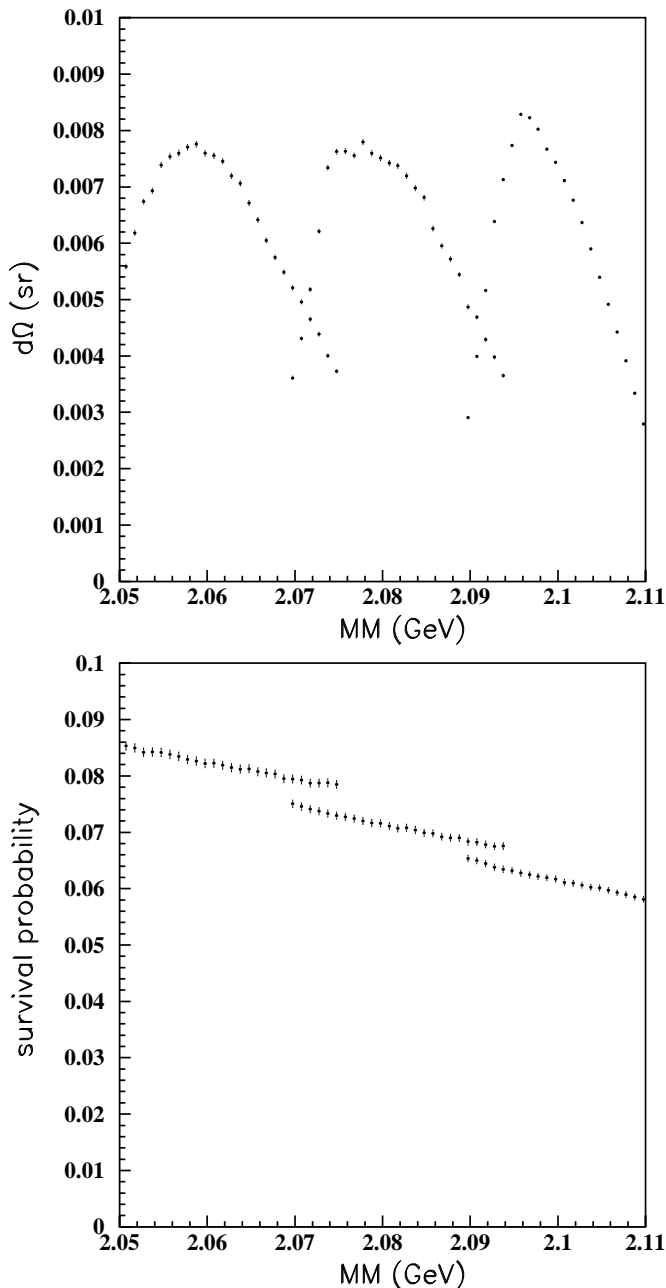


FIG. 2: Top: Acceptance correction functions. Bottom: Kaon survival probability

The kaon decay along flight path was taken into account for each individual trajectory. Path lengths were obtained from calculations of tracks in the magnetic field with the Turtle code [23]. They were checked by experimentally deduced values from time of flight measurements of pions and protons. The cross section error due to survival probability is less than 1 %. The final kaon survival probability averaged for a given missing-mass bin is presented in the bottom panel of Fig. 2.

The relative normalizations of the three different spectrograph settings were deduced from luminosity monitors located in the target area which were independent of the spectrograph settings. The relative normalization errors due to the luminosity measurement were negligibly small. The relative normalization errors due to the acceptance corrections were estimated to be less than 2 %.

The absolute cross section normalization was determined by measuring the luminosity as described in [5]. At the beginning of each beam period we calibrated two luminosity monitors counting the left and right scattered particles from the target as function of the number of beam particles. To this end, the beam current was highly reduced in order to count the beam particles individually with a fast scintillator hodoscope in the beam dump. For the calibration, the dependence of the luminosity signal on the beam intensity was fitted by a linear function. The resulting beam intensity error amounted to about 5 %. The density of the bubble-free liquid hydrogen target ($\rho = 0.0775$ g/cm³) with 1 μ m thick mylar foil windows was kept constant by stabilizing the temperature to 15.0 ± 0.5 K using a high-precision temperature control [24].

The target thickness, i.e. the distance between entrance and exit window (nominal 1 cm) was precisely measured with a calibrated optical telescope. The target thickness error was about 5 %. The overall systematic normalization error was estimated to be 10 %. The missing-mass spectrum is shown in Fig. 4.

In view of the predicted narrow dibaryon resonances in the Λp system a high missing-mass resolution was required. The missing-mass resolution depends on the spread of the beam momentum, the beam spot size at the target, the target thickness, the momentum resolution of the magnetic spectrograph and the 1 MeV bin width. The effective resolution function $f(M, M')$ has been deduced in [1] by a least-square fit to the sharply rising double differential cross section $d^2\sigma/(d\Omega_K dM_{\Lambda p})$ near the (Λp) threshold, see Fig. 4. It can be represented by a Gaussian density distribution with a standard deviation of $\sigma_M = 0.84$ MeV. It should be noted that the width σ_M depends mainly on the spread of the beam momentum and the 1 MeV bin width. The contributions from the target thickness, the beam spot size at the target (3 mm) and the momentum resolution of the magnetic spectrograph are negligibly small. Since the COSY beam has not been changed during the measurements the resolution function $f(M, M')$ is constant in the full range of missing-masses.

III. PREDICTIONS OF NARROW DIBARYON RESONANCES IN THE Λp SYSTEM

Concerning the predictions of narrow $S = -1$ dibaryon resonances in the Λp system we refer to calculations of Aerts and Dover [8, 9]. They studied the decay width of strangeness $S = -1$ six-quark bag states predicted by the Nijmegen group [10, 11] on the basis of the MIT bag model. The predicted $S = -2$ H-dibaryon with quark structure $(Q^6)_1$ is the strongest bound dibaryon state [7]. In case of $S = -1$ states the color-magnetic force is less strong and hence the lowest $(Q^6)_1$ state is predicted to be about 120 MeV above the ΛN threshold [11].

The lowest-lying $S = -1$ resonances exhibit a quark cluster structure $Q^4 \otimes Q^2$. The singlet state D_s (1P_1) is located at about 55 MeV above the Λp threshold, and the triplet states D_t ($^3P_{0,1,2}$) are located at about 95 MeV above the Λp threshold. This corresponds for the lowest-lying dibaryon D_s to an invariant mass of 2109 MeV which is located between the Λ -proton and Σ -nucleon thresholds. The resonance D_s is a singlet state of a four quark-two quark cluster configuration $(Q^4)_3 \otimes (Q^2)_{3^*}$ with angular momentum $L = 1$, spin $S = 0$, total angular momentum and parity $J^P = 1^-$. The isospin is $I = 1/2$. The only particle-decay channel is the 1P_1 wave of the Λp system. This decay is hindered by the relative P -wave centrifugal barrier between the two clusters and the stability of the $(Q^4)_3 \otimes (Q^2)_{3^*}$ configuration against dissociation into two color singlet (Q^3) clusters. In other words, the Λp channel represents only a small piece of the D_s wave function. As a consequence, the total width Γ of D_s is predicted to be less than about 1 MeV [8, 9]. In addition, the predicted width depends on the resonance mass M_r , see Fig. 2 of [8]. We note that there is a rather large uncertainty of the predicted resonance masses M_r . Therefore, the production of the dibaryon D_s has been studied theoretically in the invariant mass range between 2.06 and 2.10 GeV/ c^2 [9]. In the present paper, we consider the invariant mass range between 2.058 and 2.105 GeV/ c^2 which corresponds to resonance energies E_r between 4 MeV and 51 MeV.

We deduce from Fig. 2 of [8] $\Gamma = 26$ keV for $M_r = 2109$ MeV ($E_r = 55$ MeV). Extrapolating to $M_r = 2096.5$ MeV ($E_r = 42.5$ MeV) yields $\Gamma = 15$ keV. We note that these total widths are lower limits. The decay width calculation involves an integration of the dibaryon wave function over the radial distance r . Such integrals are dominated by the contributions from small radii r which give a lower limit of Γ . Including the larger radii gives an upper limit. Then, the lower limit has to be multiplied by 64/3 to obtain an upper limit. In addition, the width increases by about a factor of 2 if Hulthén rather than oscillator wave functions are used. Thus, an upper limit is obtained if the lower limit is multiplied by a factor 128/3 [8]. We get the following ranges for the predicted total width: $\Gamma = [15, 640]$ keV for $M_r = 2096.5$ MeV ($E_r = 42.5$ MeV) and $\Gamma = [26, 1109]$ keV for $M_r = 2109$ MeV ($E_r = 55$ MeV). As one can see, the

predicted widths depend strongly on the invariant mass of the dibaryon.

We emphasize that a possible resonance D_s below the $\Sigma^+ n$ threshold (2128.935 MeV) can only decay into $\Lambda + p$. Therefore, the elastic width Γ_{el} can be assumed to be equal to the total width Γ , i.e. $\Gamma_{el}/\Gamma = 1$.

IV. SIMULATION OF A NARROW RESONANCE IN THE TOTAL CROSS SECTION OF $\Lambda p \rightarrow \Lambda p$ SCATTERING

We first discuss the effect of a single isolated narrow resonance in the total cross section of the $\Lambda + p \rightarrow \Lambda + p$ scattering and we compare it with existing data from bubble chamber experiments [25–30]. We start with the ansatz

$$\sigma_{tot} = \sigma_{tot}^{nr} + \sigma_{tot}^r. \quad (1)$$

Here, σ_{tot}^{nr} is the nonresonant part and σ_{tot}^r the resonant part of the total cross section. Assuming S -waves and taking the effective range approximation the nonresonant part of the total cross section may be written [31]

$$\sigma_{tot}^{nr} = \frac{1}{4} \frac{4\pi}{k^2 + \left(-\frac{1}{a_s} + \frac{r_s k^2}{2}\right)^2} + \frac{3}{4} \frac{4\pi}{k^2 + \left(-\frac{1}{a_t} + \frac{r_t k^2}{2}\right)^2} \quad (2)$$

Here, a_s and a_t are the singlet and triplet scattering lengths, r_s and r_t are the singlet and triplet effective ranges and $k = q/\hbar$ is the wave number corresponding to the c.m. momentum q . We take the effective range parameters $a_s = -2.43$ fm, $r_s = 2.21$ fm, $a_t = -1.56$ fm, $r_t = 3.7$ fm as determined in [1].

The resonant part of the total cross section can be approximated by the Breit-Wigner resonance formula,

$$\sigma_{tot}^r = \frac{2J_r + 1}{(2s_1 + 1)(2s_2 + 1)} \frac{\pi}{k^2} \frac{\Gamma^2}{(E - E_r)^2 + (\Gamma/2)^2}. \quad (3)$$

Here, E is the total kinetic energy in the c.m. system, $k = q/\hbar$ is the wave number corresponding to the c.m. momentum q and we assume $\Gamma_{el} = \Gamma$. We take the narrow resonance D_s as an example. The resonance is predicted to occur in the partial wave 1P_1 . Therefore, we assume $J_r = 1$. As resonance energy we take $E_r = 42.5$ MeV corresponding to $\sqrt{s} = 2096.5$ MeV of the peak seen by Siebert et al. [21]. At this energy, the nonresonant potential scattering is dominated by the partial waves 1S_0 and 3S_1 and the nonresonant contributions of the P waves can be neglected in Eq. (2). The same holds true for a possible interference between the resonant and nonresonant 1P_1 amplitude in Eq. (3). The Jülich hyperon-nucleon model [32] yields for instance phase shift predictions with $\delta < 5^\circ$ for 1P_1 and $^3P_{0,1,2}$.

In Fig. 3 we show the effect of the D_s resonance with $J_r = 1$ and $E_r = 42.5$ MeV for $\Gamma = 500$ keV without taking the effect of a finite energy resolution into

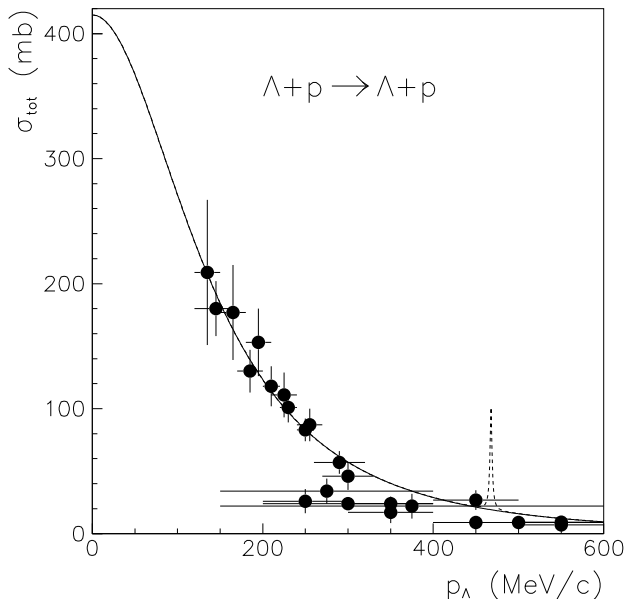


FIG. 3: Total $\Lambda p \rightarrow \Lambda p$ cross section vs laboratory momentum p_Λ . Solid line: effective range approximation of σ_{tot}^{nr} . Dashed line: Simulation of a resonance signal σ_{tot}^r without folding with the effective resolution function for a resonance in the 1P_1 channel with $E_r = 42.5$ MeV and $\Gamma = 500$ keV. The data are from [25–30]

account. We take $\Gamma = 500$ keV as an example. A resonance with $\Gamma = 15$ keV would appear as a sharp needle in Fig. 3. The maximum of the resonance signal at $E = E_r$ is very large, $\sigma_{tot}^r(E_r)/\sigma_{tot}^{nr}(E_r) = 4.2$. Such a resonance could be easily observed provided the effective energy resolution would be sufficiently high. Unfortunately, the bubble chamber experiments suffer from low-energy resolution as well as from low statistical accuracy. For instance data points near the resonance exhibit bin widths of 50 MeV/c ($\Delta E_{cm} = 8.4$ MeV) and 100 MeV/c ($\Delta E_{cm} = 16.8$ MeV). These bin widths are so large that the effect of a narrow resonance is completely diluted.

V. NARROW RESONANCE EMBEDDED IN THE CONTINUUM OF $pp \rightarrow K^+(\Lambda p)$

The reaction $pp \rightarrow K^+(\Lambda p)$ can be described by factorizing the reaction amplitude in terms of a production amplitude and a final state enhancement factor. The method of parametrizing the FSI enhancement factor by the inverse Jost function [31] is described in [1, 33]. Taking the spin statistical weights into account the nonreso-

nant double differential cross section may be written as

$$\frac{d^2\sigma^{nr}}{d\Omega_K dM_{\Lambda p}} = \Phi_3 \left[\frac{1}{4} |M_s|^2 \frac{q^2 + \beta_s^2}{q^2 + \alpha_s^2} + \frac{3}{4} |M_t|^2 \frac{q^2 + \beta_t^2}{q^2 + \alpha_t^2} \right]. \quad (4)$$

Here, $|M_s|^2$ and $|M_t|^2$ are the singlet and triplet production matrix elements squared, q the internal c.m.-momentum of the Λp system, $\alpha_s, \beta_s, \alpha_t, \beta_t$ the singlet and triplet potential parameters, and Φ_3 the ratio of the three-body phase space distribution and the incident flux factor. The potential parameters α and β can be used to establish phase-equivalent Bargmann potentials [34]. They are related to the scattering lengths a , and effective ranges r of the low-energy S wave scattering, $\alpha = (1 - \sqrt{1 - 2r/a})/r$, $\beta = (1 + \sqrt{1 - 2r/a})/r$. Thus, the shape of the missing-mass spectrum (4) depends on $|M_s|^2, |M_t|^2$ and the effective range parameters a_s, r_s, a_t and r_t . These parameters have been deduced by a combined fit of the $pp \rightarrow K^+(\Lambda p)$ data and the total Λp cross section data [1]. In principle one could equally well take a fit with three spin-averaged parameters $|\bar{M}|^2, \bar{a}$ and \bar{r} (see [1]) or another three-parameter description of the missing-mass spectrum as in e.g. [35].

A narrow resonance D_s embedded in the Λp continuum of the reaction $pp \rightarrow K^+(\Lambda p)$ can be described using the following equation

$$\frac{d^2\sigma^r}{d\Omega_K dM_{\Lambda p}} = \frac{d\sigma^r}{d\Omega_K} \frac{2M_{\Lambda p}}{\pi} \frac{M_r \Gamma}{(M_{\Lambda p}^2 - M_r^2)^2 + M_r^2 \Gamma^2}. \quad (5)$$

This equation corresponds to a formula derived by Pilkuhn [36] in order to describe the reaction $a + b \rightarrow c + d$ where d is a narrow resonance decaying into particles 1 and 2, $d \rightarrow 1 + 2$. Here, $d\sigma^r/d\Omega_K$ represents the differential cross section for the production of the resonance D_s , M_r the resonance mass, Γ the total width and $M_{\Lambda p}$ the invariant mass of the Λp system. The relativistic Breit-Wigner form in (5) can be approximated by the corresponding nonrelativistic Breit-Wigner form,

$$\frac{d^2\sigma^r}{d\Omega_K dM_{\Lambda p}} = \frac{d\sigma^r}{d\Omega_K} \frac{1}{2\pi} \frac{\Gamma}{(M_{\Lambda p} - M_r)^2 + (\Gamma/2)^2}. \quad (6)$$

In passing, we note that the integral over the Breit-Wigner distribution is normalized to one,

$$\int_{-\infty}^{\infty} \frac{1}{2\pi} \frac{\Gamma}{(M_{\Lambda p} - M_r)^2 + (\Gamma/2)^2} dM_{\Lambda p} = 1. \quad (7)$$

The double differential cross section including a possible narrow resonance D_s in the Λp system may be written by combining (4) and (5),

$$\frac{d^2\sigma}{d\Omega_K dM_{\Lambda p}} = \frac{d^2\sigma^{nr}}{d\Omega_K dM_{\Lambda p}} + \frac{d^2\sigma^r}{d\Omega_K dM_{\Lambda p}}. \quad (8)$$

We mention again that possible P wave contributions in the Λp interaction are negligibly small at low Λp c.m. energies [32]. Therefore, the effect of an interference with

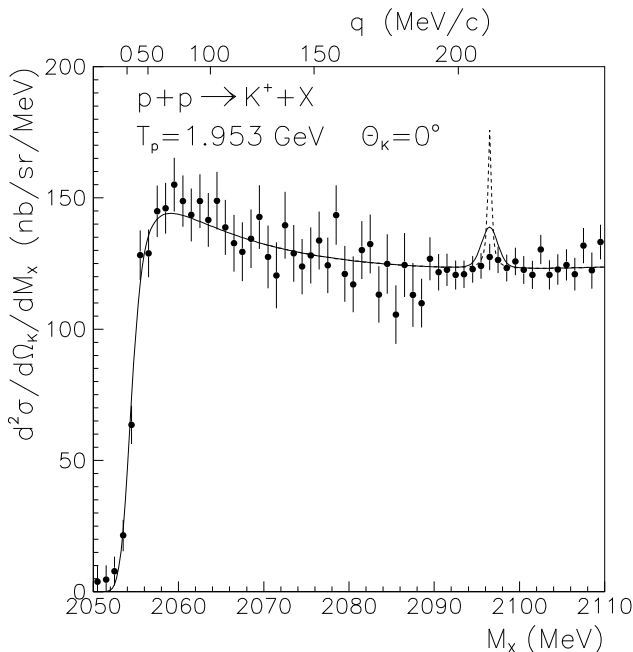


FIG. 4: Missing-mass spectrum of the reaction $p + p \rightarrow K^+ + X$ with $X = (\Lambda p)$ measured at $T_p = 1.953$ GeV and $\Theta_K = 0^\circ$. The upper axis indicates the c.m. momentum q of the Λp system. Solid line: FSI fit curve with resonance signal excluded by the χ^2 test. Dashed line: Same resonance signal without folding with the effective resolution function. Resonance parameters: $d\sigma^r/d\Omega_K = 42$ nb/sr, $M_r = 2096.5$ MeV, $\Gamma = 500$ keV.

the nonresonant 1P_1 decay channel can be neglected. The relative contribution of a narrow resonance D_s to the double differential cross section depends on the production cross section $d\sigma^r/d\Omega_K$ for $(pp \rightarrow K^+ D_s)$. In the present paper, we deduce upper limits for the production cross section.

VI. UPPER LIMITS FOR A NARROW RESONANCE

A narrow $S = -1$ resonance D_s as predicted by [8, 9] is not visible in the data (see Fig. 4). There may be two reasons for it. (i) The predicted narrow $S = -1$ resonance D_s does not exist at all in the invariant mass region below 2110 MeV. (ii) The production cross section $d\sigma^r/d\Omega_K(pp \rightarrow K^+ D_s)$ is too small. In the following, we deduce upper limits for the production cross section in the given invariant mass range.

A. Effective resolution function

We note that the theoretical expression for the double differential cross section (8) must be folded with the effective

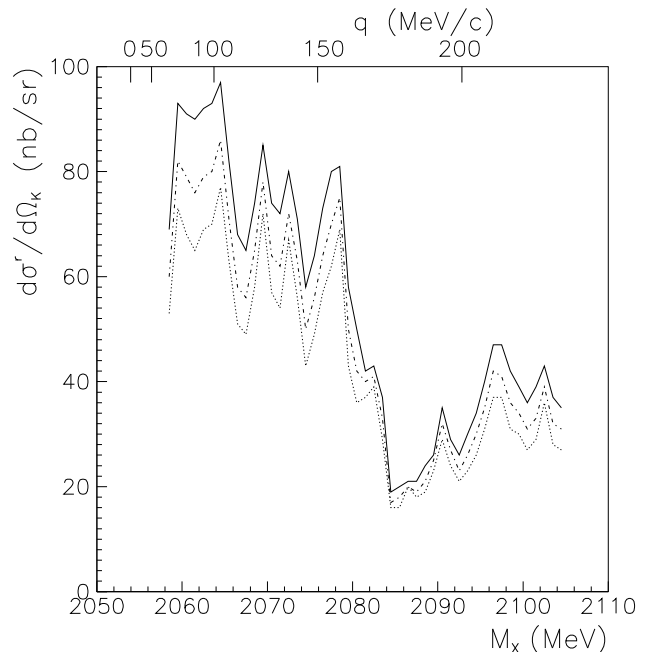


FIG. 5: Upper limits with 99% confidence level of the production cross section $d\sigma^r/d\Omega_K$ (nb/sr) for $\Gamma = 100$ keV (dotted line), $\Gamma = 500$ keV (dashed-dotted line), and $\Gamma = 1.0$ MeV (solid line).

resolution function $f(M, M')$ before comparing with the data. The effective resolution function is represented by a Gaussian density distribution with a standard deviation of $\sigma_M = 0.84$ MeV. It was deduced in [1] by a least-square fit to the sharply rising double differential cross section $d^2\sigma/(d\Omega_K dM_{\Lambda p})$ near the (Λp) threshold, see Fig. 4. As discussed in Sec. II the effective resolution function is constant in the full missing-mass range between 2050 MeV and 2110 MeV.

B. χ^2 test

The compatibility of the experimental data with the hypothesis of a narrow resonance in the invariant mass region below 2110 MeV has been studied using a χ^2 test. Using this method we determine the upper limits of the production cross section $d\sigma^r/d\Omega_K(pp \rightarrow K^+ D_s)$ in the Λp invariant mass range between 2058 and 2105 MeV, see Fig. 5. We only include the data measured within ± 4 MeV around the assumed resonance mass M_r . The data are binned in 1 MeV wide bins; thus nine data points are used for the χ^2 test. Including more data would dilute the information, since an excursion with a width of about $\sigma_M = 0.84$ MeV is negligibly small for $|M_{\Lambda p} - M_r| \geq 5$ MeV, see Fig. 4.

The χ^2 obtained from a comparison of the hypothesis to the data is subject to a standard χ^2 test [37]. We

assume that the χ^2 is statistically distributed according to a χ^2 probability density function $f(\chi^2, n_d)$ with the appropriate number of degrees of freedom, n_d . Thus, we can calculate the confidence level (CL) with which we can support or falsify the hypothesis. If the hypothesis is true, the probability P to obtain a larger χ^2 than the observed one in an infinite repetition of the experiment is given by

$$P = \int_{\chi^2}^{\infty} f(z, n_d) dz. \quad (9)$$

If P is very small, either the hypothesis is wrong or the current measurement suffers from a very unlikely statistical fluctuation. The CL for *excluding* the hypothesis of a narrow resonance with a certain production cross section is given by $1 - P$. Thus, we have $P = 1\%$ for a confidence level of 99 %, i.e. the probability to observe a larger χ^2 when repeating the experiment is only 1 %.

We use the number of data points within the invariant mass interval $M_r \pm 4$ MeV as the number of degrees of freedom, n_d . Thus, we neglect the degrees of freedom introduced by fitting the Λp missing-mass spectrum. Three fit parameters, i.e. the spin-averaged parameters \bar{M}^2 , \bar{a} and \bar{r} , are sufficient in order to reproduce the measured missing-mass spectrum. They cannot be considered as free parameters in the present hypothesis test. We mention that there is no significant effect on the deduced upper limits due to the errors of those parameters. Thus, we have $n_d = 9$ and $\chi^2/n_d = 2.41$ for $P = 1\%$.

C. Upper limits for the production cross section $d\sigma^r/d\Omega_K(pp \rightarrow K^+ D_s)$

The upper limits for the production cross section $d\sigma^r/d\Omega_K(pp \rightarrow K^+ D_s)$ are shown in Fig. 5 for $\Gamma = 100$ keV (dotted line), $\Gamma = 500$ keV (dashed-dotted line), and $\Gamma = 1.0$ MeV (solid line). We note that the upper limits depend slightly on the total width Γ . For very small total widths, i.e. $\Gamma \leq 100$ keV one can take the upper limits for $\Gamma = 100$ keV. The differences between the upper limits for $\Gamma < 100$ keV and $\Gamma = 100$ keV are negligible.

The upper limits exhibit rather strong fluctuations due to the statistical fluctuations of the data. In the lower part of the missing-mass spectrum between 2058 and 2080 MeV, the upper limit varies between 43 and 77 nb/sr for $\Gamma = 100$ keV, 50 and 86 nb/sr for $\Gamma = 500$ keV, and 50 and 97 nb/sr for $\Gamma = 1.0$ MeV. In the region between 2080 and 2105 MeV the upper limits are roughly a factor of 2 smaller. This is because four data points between 2080 and 2090 MeV are below the fit curve and that the missing-mass region between 2090 and 2105 MeV has been measured with a much higher statistical accuracy. There, the upper limit varies between 21 and 37 nb/sr for $\Gamma = 100$ keV, 23 and 42 nb/sr for $\Gamma = 500$ keV and 26 and 47 nb/sr for $\Gamma = 1.0$ MeV.

VII. CONCLUSION AND DISCUSSION

The effect of a single isolated narrow resonance in the total cross section of the free Λp scattering has been studied (see Fig. 3). A narrow strangeness $S = -1$ resonance in the 1P_1 channel is predicted to be characterized by a rather large resonance signal. Therefore, the free Λp scattering would be ideal to search for a narrow resonance. However, the effective resolution and the statistical accuracy of existing Λp scattering data are too low for a stringent upper limit search.

The reaction $pp \rightarrow K^+ + (\Lambda p)$ has been measured at $T_p = 1.953$ GeV and $\Theta = 0^\circ$ with a high missing-mass resolution. A narrow strangeness $S = -1$ resonance D_s as predicted by [8, 9] is not visible in the data (see Fig. 4). The missing-mass range between 2090 and 2110 MeV has been studied with an especially high statistical accuracy. The apparent small enhancement near 2096.5 MeV (see Fig. 4) corresponds to an excursion of only 1.5 standard deviations. Upper limits of the production cross section $d\sigma^r/d\Omega_K(pp \rightarrow K^+ D_s)$ are deduced for resonance energies between 2058 and 2105 MeV, see Fig. 5. In the missing-mass region between 2090 and 2105 MeV the upper limit varies between 21 and 37 nb/sr for $\Gamma = 100$ keV, 23 and 42 nb/sr for $\Gamma = 500$ keV and 26 and 47 nb/sr for $\Gamma = 1.0$ MeV.

Unfortunately, there are no theoretical predictions of the production cross section $d\sigma^r/d\Omega_K(pp \rightarrow K^+ D_s)$. Aerts and Dover gave estimates of production cross sections of strangeness $S = -1$ dibaryon states $D_{s,t}$ for the reactions $K^- + d \rightarrow \pi^- + D_t$ and $K^- + {}^3\text{He} \rightarrow \pi^- + n + D_s$ [9]. Differential cross sections of order 1 $\mu\text{b/sr}$ were predicted.

The transition $pp \rightarrow K^+ D_s$ from the pp entrance channel requires orbital angular momentum transfer $\Delta l = 0$ and 2. Thus, the production cross section $d\sigma^r/d\Omega_K(pp \rightarrow K^+ D_s)$ should peak at the scattering angle $\Theta = 0^\circ$. The reaction $pp \rightarrow K^+ (\Lambda p)$ has been measured at $\Theta = 0^\circ$, i.e. at the maximum of the angular distribution. Therefore, one expects maximum sensitivity when searching for a narrow strangeness $S = -1$ resonance D_s with the reaction $pp \rightarrow K^+ D_s$.

It is interesting to note that the upper limit of the production cross section $d\sigma^r/d\Omega_K(pp \rightarrow K^+ D_s)$ is nearly independent of the total width Γ as long as $\Gamma < \sigma_M$. This is due to the fact that the integral over the Breit-Wigner distribution (6) is constant (equal to one), that means it does not depend on Γ . Thus, folding a narrow Breit-Wigner distribution (6) with a wider Gaussian distribution yields a resonance signal which depends only on the production cross section and the effective resolution width σ_M but not on Γ . This is a great advantage when searching for an extremely narrow resonance. The sensitivity in the search for a narrow resonance can be increased substantially by increasing the missing-mass resolution, i.e. by decreasing the resolution width σ_M . Another less efficient way is to increase the statistical and systematic accuracy.

In contrast to previous experiments, the present experiment exhibits a rather high missing-mass resolution. In previous experiments, the missing-mass resolution has been very much lower. Hogan et al. [14] studied the reaction $pp \rightarrow K^+(\Lambda p)$ at bombarding energies between 2.5 and 3.0 GeV and scattering angles 20° , 30° and 40° . The momentum resolution of the detected kaons ($\Delta p/p$) ranged from 1.5 % at 20° to 3 % at 40° . The number of data points per MeV/ c and the effective missing-mass resolution was so low that they were not able to see the marked FSI enhancement near the Λp threshold. Reed et al. [13] studied the reaction $pp \rightarrow K^+(\Lambda p)$ at 2.40 and 2.85 GeV and 0° , 17° and 32° . The momentum resolution of the outgoing kaons at 0° corresponded to $\Delta p/p = 1.5$ %. This resolution was also too low to search for a narrow resonance with $\Gamma \leq 1$ MeV. The same holds true for the exclusive measurements of the reaction $pp \rightarrow K^+\Lambda p$ by the COSY-TOF Collaboration [15–17].

As mentioned in the introduction, Siebert et al. [21] performed the first high-resolution study of the reaction $pp \rightarrow K^+(\Lambda p)$. Sharp peaks in the missing-mass spectra seen there have not been confirmed by the present experiment. In this context we note that the reaction

$pp \rightarrow K^+(\Lambda p)$ has been studied [21] at $T_p = 2300$ MeV at four scattering angles, 6° , 8.3° , 10.3° , and 12° . The peak near 2096.5 ± 1.5 appeared only in the missing-mass spectrum measured at 10.3° but not at 6° , 8.3° and 12° . At $T_p = 2700$ MeV, the interesting missing-mass range has been studied at two scattering angles, 12.6° and 20° . The peak near 2098.0 ± 1.5 appeared only at $\Theta_K = 12.6^\circ$ but not at 20° . We therefore conclude that the peaks near 2096.5 ± 1.5 and 2098.0 ± 1.5 MeV must be attributed to a statistical fluctuation of the nonresonant cross section.

Summarizing, the reaction $pp \rightarrow K^+(\Lambda p)$ has been measured with a high missing-mass resolution at $T_p = 1.953$ GeV and $\Theta = 0^\circ$. A narrow strangeness $S = -1$ resonance D_s is not visible in the missing-mass spectrum. Upper limits for the production cross section of $pp \rightarrow K^+D_s$ have been deduced.

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